

# Integrodifference equations in spatial ecology

## Lecture 8: Spread with Allee effect

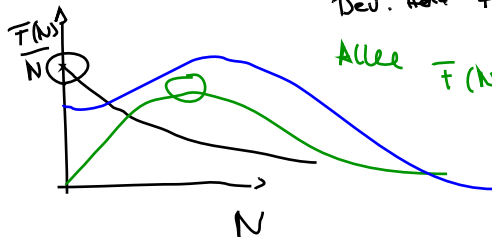
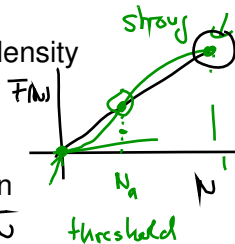
Frithjof Lutscher

# Allee effect

$$N_{t+1} = \bar{F}(N_t)$$

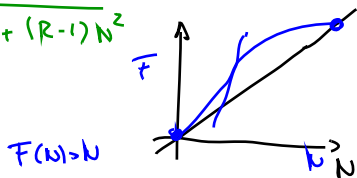
→ per-capita growth rate maximal at intermediate density

- strong vs weak
- examples
- mathematical difficulty: bistability, no linearization



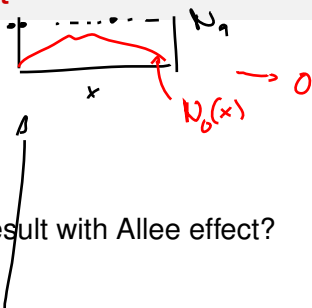
Dev. hold  $F(N) = \frac{RN^2}{1+(R-1)N}$

Allee  $\bar{F}(N) = \frac{RN^2}{1+(R-1)N^2}$



## Spread with Allee effect

- Speed of spread?
- Spread or not?
- Traveling waves?
- Recall the steady-state result with Allee effect?

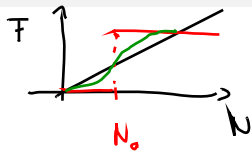


Idea: If  $\sigma^2 \ll 1$  then we get persistence in the IDE on  $\mathbb{R}^d$  domain if  $F(1)=1$  is stable for non-spatial model

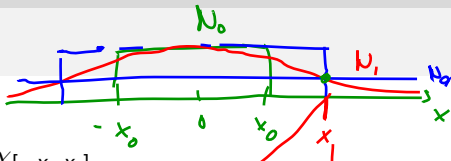
Answer: No

# A caricature Allee function (Kot et al 1996)

$$F(N) = \begin{cases} 0, & N < N_a \\ 1, & N \geq N_a \end{cases}$$



# Calculating spatial extent I



$$N_0(x) = \chi_{[-x_0, x_0]}$$

$$\underline{F(N_0)} = N_0$$

$$N_1(x) = \int_{-\infty}^{\infty} k(x-y) F(N_0(x,y)) dy = \int_{-\infty}^{\infty} k(x-y) \chi_{[-x_0, x_0]} dy$$

$$= \int_{-x_0}^{x_0} k(x-y) dy$$

$$\int_{x-x_0}^{x+x_0} k(z) dz$$

$$\int_{x_1-x_0}^{x_1+x_0} k(z) dz = N_1$$

$$y = -x_0 \quad z = x+x_0$$

$\Rightarrow x_1$  is determined by

$$F(N_1(x)) = \chi_{[-x_1, x_1]}$$

## Calculating spatial extent II

$$\int_{x_{t+1}-x_t}^{x_{t+1}+x_t} K(y) dy = N_a$$

$= 2x_0$  (pointing to the upper limit)

$= 0$  (pointing to the lower limit)

This is a 1-D iteration

Case 1:  $x_1 < x_0$  in general  $x_{t+1} < x_t$

Case 2:  $x_1 > x_0$  in general  $x_{t+1} > x_t$

threshold initial size:

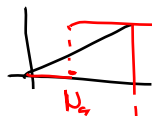
$$x_{t+1} = x_t = x_c$$

Critical initial size

Requires  $N_a < 1/2$ .

$$\int_0^{2x_c} K(y) dy = N_a$$

$< 1/2$  (pointing to the integral result)





# Spatial extent and Laplace kernel

Defining eq.

$$\int_{x_{t+1}-x_t}^{x_{t+1}+x_t} \frac{a}{2} e^{-a|y|} dy = N_a$$

Critical extent

$$\int_0^{2x_c} \frac{a}{2} e^{-ay} dy = -\frac{1}{2} e^{-ay} \Big|_0^{2x_c} = \frac{1}{2} (1 - e^{-2ax_c}) = N_a$$

$$1 - e^{-2ax_c} = 2N_a$$

$$1 - 2N_a = e^{-2ax_c}$$

$$-2ax_c = \ln(1 - 2N_a)$$

$$x_c = \frac{-1}{2a} \ln(1 - 2N_a)$$

$$N_a < \frac{1}{2}$$



# Iterations for $x_{t+1} > x_t$

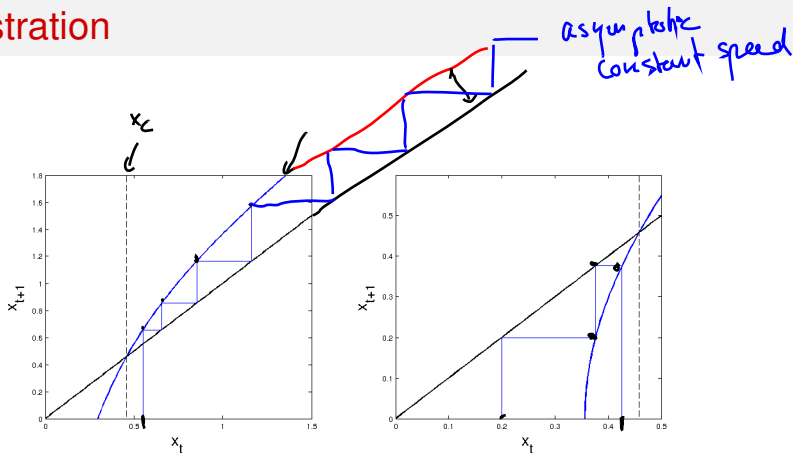
$$\int_{x_{t+1} - x_t}^{x_{t+1} + x_t} \frac{a}{2} e^{-ay} dy = N_a \quad \text{solve for } x_t$$

$$= -\frac{1}{2} e^{-ay} \Big|_{x_{t+1} - x_t}^{x_{t+1} + x_t} = -\frac{1}{2} \left( e^{-a(x_{t+1} + x_t)} - e^{-a(x_{t+1} - x_t)} \right)$$

$$= + \frac{1}{2} e^{-a(x_{t+1})} \left( e^{+ax_t} - e^{-ax_t} \right) = + e^{-ax_{t+1}} \sinh(ax_t)$$

$$e^{-ax_{t+1}} \sinh(ax_t) = N_a \quad \Rightarrow \quad x_{t+1} = \frac{1}{a} \ln \left( \frac{\sinh(ax_t)}{N_a} \right)$$

# Illustration



$$x_{t+1} = \frac{1}{a} \ln \left( \frac{\sinh(ax_t)}{N_a} \right)$$

$$\sinh(ax_t) \approx \frac{1}{2} e^{ax_t} \quad \text{für } x_t \text{ large}$$

$$x_{t+1} \approx \frac{1}{a} \ln \left( \frac{e^{ax_t}}{2N_a} \right) = x_t - \frac{\ln(2N_a)}{a} \quad N_a < \frac{1}{2}$$

# The asymptotic spreading speed

$\Gamma W:$



$$N_a = \int_{-\infty, 0}$$

$$F(N_a) = \int_{-\infty, c}$$

Explicit, finite for ~~all~~ many kernels, including Cauchy **but!**

$$\int_0^{c^*} K(z) dz = \frac{1}{2} - N_a$$

$$x_{t+n} - x_t = c$$

$$x_{t+n} + x_t \rightarrow \infty \text{ as } t \rightarrow \infty$$

$$\int_{x_{t+n}-c}^{x_{t+n}+c} K dy = N_a$$

**but!**

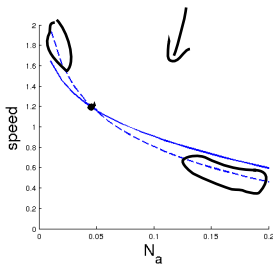
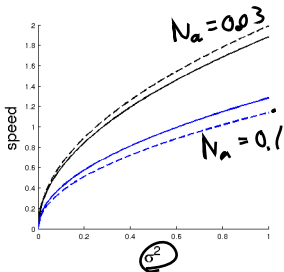
$$\int_c^{\infty} K dy = N_a$$

$$\int_{-\infty}^c K dy = -N_a$$

$$\int_{-\infty}^{\infty} K dy = \frac{1}{2}$$

$$\int_c^{\infty} K dy = \frac{1}{2}$$

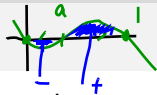
$$\int_c^{\infty} K dy = \int_c^{\infty} K dy + \frac{1}{2}$$



Gauss (solid) vs Laplace (dashed)

$$N_a = \int_c^{\infty} K dy = \int_c^{\infty} K dy + \frac{1}{2}$$

# Direction of the wave



The reaction-diffusion case with bistable nonlinearity

Nagumo

$$0 < a < 1$$

bistability

$$u_t = \underbrace{D u_{xx}}_{\text{random}} + \underbrace{f(u)}_{\text{growth}}$$

$$f(u) = u(1-u)(u-a)$$



FW:  $u(x-ct) = u(x,t) = u(z)$

$$-cu' = Du'' + f(u) \rightarrow Du'' + cu' + f(u) = 0$$

Multiply by  $u'$   $Du''u' + cu'^2 + f(u)u' = 0$

sign(c) =  
sign  $\int f(u) du$   
 $\uparrow$   
0

integrate:  $D \int_{-\infty}^{\infty} u'' u' dz + c \int_{-\infty}^{\infty} u'^2 dz + \int_{-\infty}^{\infty} f(u) u' dz = 0$

$$\rightarrow \frac{(u')^2}{2} \Big|_{-\infty}^{\infty} = 0$$

$$= \int_0^1 f(u) du$$

$$\implies c = \frac{\int_0^1 f(u) du}{\int u'^2 dz}$$

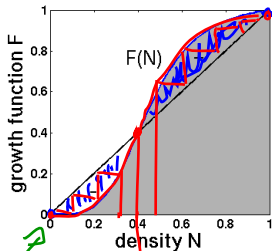
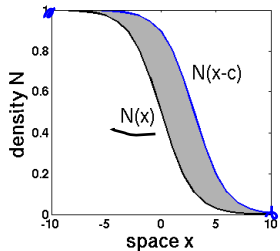
# The integrodifference case - visual



$$N_{t+1}(x) = \int_{-\infty}^{\infty} k(x-y) \underbrace{F(N_t(y)) dy}_{\text{Allee strong}}$$

for  $\tilde{u} = f(u)$

here:  $N_{t+1} = F(N_t)$



$$\int_0^1 (F(N) - N) dN$$

1 = grey + white =  $\int_0^1 F(N) dN + \int_0^1 F^{-1}(y) dy$

# The theorem (Wang et al 2002) |

*Kol, Neubert*

Consider the IDE  $N_{t+1}(x) = (K * F(N_t))(x)$  with monotone growth function  $F$  and steady states  $N = 1$  and  $N = 0$ . Assume that there is a monotone decreasing travelling front with with speed  $c$  and profile  $N(x)$ , connecting the two states.

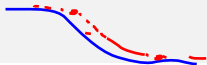
Furthermore, assume that  $F$  and  $N$  are real analytic functions, that derivatives of any order of  $N$  vanish as  $x \rightarrow \pm\infty$  and that the derivatives  $d^i F(N(x))/dx^i$  are bounded uniformly in  $i$ .

Then we have the following relation for the sign of the speed of the travelling front:

$$\text{sign}(c) = \text{sign} \left( \int_0^1 [F(N) - N] dN \right).$$



# Sketch of the proof



If  $N$  is decreasing then

$$c > 0 \iff N(x) < N(x-c) \quad \forall x$$

$$\frac{d}{dx} F(N(x)) = \underbrace{\frac{dF}{dN}}_{>0} \cdot \underbrace{\frac{dN}{dx}}_{<0} < 0$$

$$\iff (N(x) - N(x-c)) \frac{dF(N(x))}{dx} > 0 \quad \forall x$$

Then  $0 < \int_{-\infty}^{\infty} (N(x) - N(x-c)) \frac{dF}{dx} dx \stackrel{\text{Lemma}}{=} \int_{-\infty}^{\infty} (N - F(N)) \frac{dF}{dx} dx$

$$y = F(N(x))$$

$$\downarrow$$

$$F^{-1}(y) = N$$

$$\int_0^1 [F^{-1}(y) - y] dy = \int_0^1 (y - F^{-1}(y)) dy$$

$$\int_0^1 (F(N) - N) dN > 0$$

$$= \int_0^1 y dy + \int_0^1 F(N) dN - 1 = -\int_0^1 N dN + 1 + \int_0^1 F dN - 1$$

# General theory (Lui 1983)

Consider the IDE  $\underline{N_{t+1}(x) = (K * F(N_t))(x)}$  where

- i)  $K$  continuous, symmetric;  $M(s)$  exists for all  $s$
- ii)  $\int_x^\infty K(y)dy \leq CK(x)$  for large  $x$ .
- iii)  $\underline{F}$  is  $C^1$  and  $F(0) = 0 = F(1) - 1 = F(N_a) - N_a$ .
- > iv)  $F'(N) \leq F'(N_a)$  for  $N \in [0, 1]$ .
- > v)  $\underline{F'(0)N} \leq F(N) \leq \underline{F'(1)(N - 1)} + 1$  for  $N \in [0, 1]$ .





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- v)  $F'(0)N \leq F(N) \leq F'(1)(N - 1) + 1$  for  $N \in [0, 1]$ .

1 There exists an asymptotic spreading speed  $c^*$  is the following sense. If  $N_0(x) = 0$  for  $x > 0$  and  $N_0(x) > N_a$  as  $x \rightarrow -\infty$ , then

$$\limsup_{t \rightarrow \infty} \max_{x > (c^* + \epsilon)t} N_t(x) = 0 \quad \text{and} \quad \liminf_{t \rightarrow \infty} \min_{x < (c^* - \epsilon)t} N_t(x) = 1.$$



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- 2 A monotone travelling wave can exist for at most one speed.

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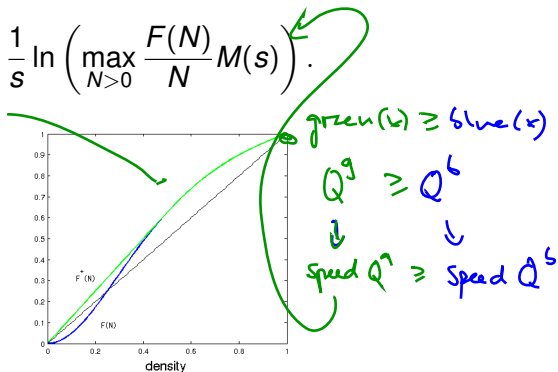
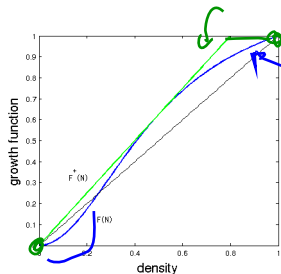
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- 2 A monotone travelling wave can exist for at most one speed.
- 3 There exists a speed  $c^* \in \mathbb{R}$  and a family of monotone travelling waves with speed  $c^*$ .

# A bound for the spreading speed

Use the monostable theory to bound the spreading speed in the bistable case

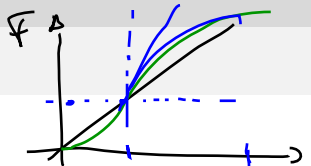
$$c^* \leq \max_{s>0} \frac{1}{s} \ln \left( \max_{N>0} \frac{F(N)}{N} M(s) \right).$$



## Another traveling wave

Assume  $F$  is as in Lui's theorem. Define

$$\underline{G(N) = F(N + N_a) - N_a}, \quad \underline{N \in [0, 1 - N_a]}.$$



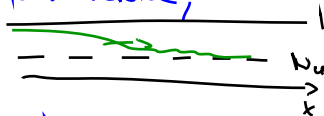
Then  $G$  is a monostable function and the IDE

$$N_{t+1}(x) = (K * G(N_t))(x)$$

has a linearly determined spreading speed and monotone traveling waves.

$$G(0) = 0, \quad G'(0) = 1 - N_a, \quad G \text{ is monotone,}$$

$$G(N) \leq G'(0) \cdot N$$



$\Rightarrow$  These are TW for  $Q^G$  from 0 to  $1 - N_a$   
 $\Rightarrow$  - - - for  $Q^F$  from  $N_a$  to 1